

فصل 7

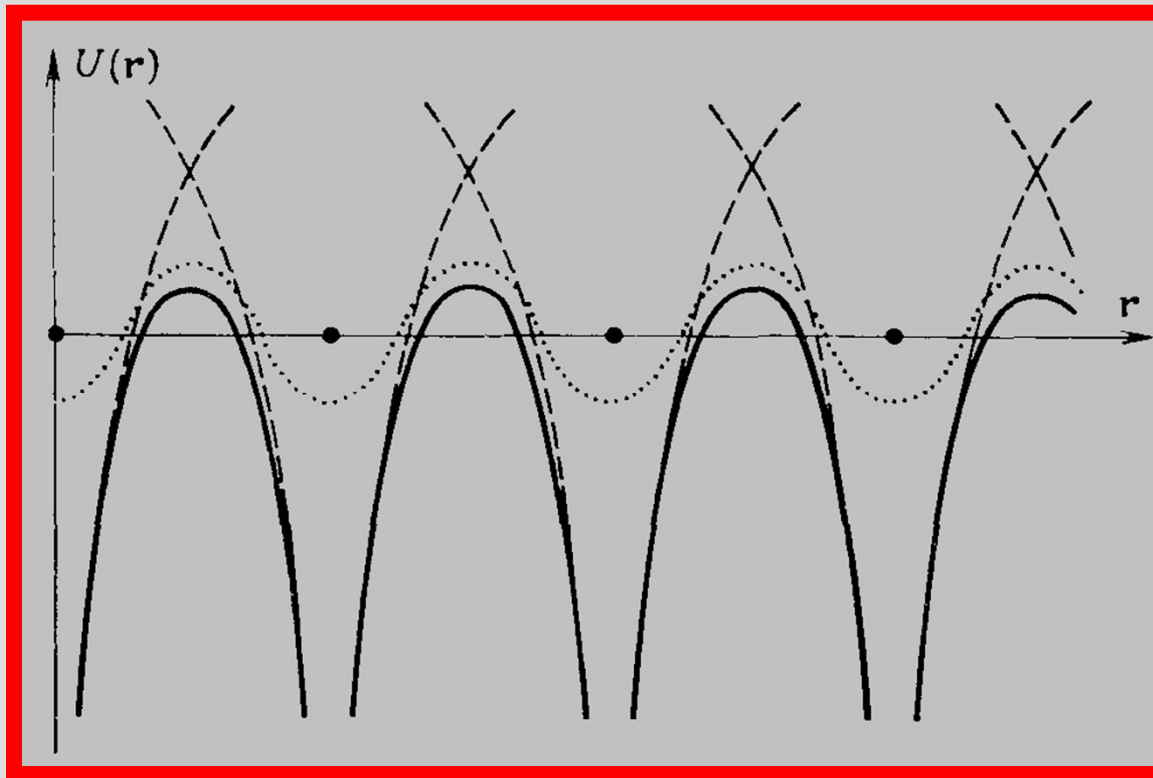
ترازهای الکترونی در یک پتانسیل تناوبی

قضیه بلوخ

ساختار نواری

مدل گرونینگ پنی

THE PERIODIC POTENTIAL

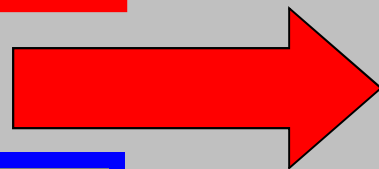


$$H\psi = \left(-\frac{\hbar^2}{2m} \nabla^2 + U(\mathbf{r}) \right) \psi = \mathcal{E}\psi$$

BLOCH'S THEOREM

Theorem.¹ The eigenstates ψ of the one-electron Hamiltonian $H = -\hbar^2 \nabla^2 / 2m + U(\mathbf{r})$, where $U(\mathbf{r} + \mathbf{R}) = U(\mathbf{r})$ for all $\mathbf{R}$ in a Bravais lattice, can be chosen to have the form of a plane wave times a function with the periodicity of the Bravais lattice:

$$\psi_{n\mathbf{k}}(\mathbf{r}) = e^{i\mathbf{k} \cdot \mathbf{r}} u_{n\mathbf{k}}(\mathbf{r})$$



$$\psi_{n\mathbf{k}}(\mathbf{r} + \mathbf{R}) = e^{i\mathbf{k} \cdot \mathbf{R}} \psi_{n\mathbf{k}}(\mathbf{r})$$

$$u_{n\mathbf{k}}(\mathbf{r} + \mathbf{R}) = u_{n\mathbf{k}}(\mathbf{r})$$

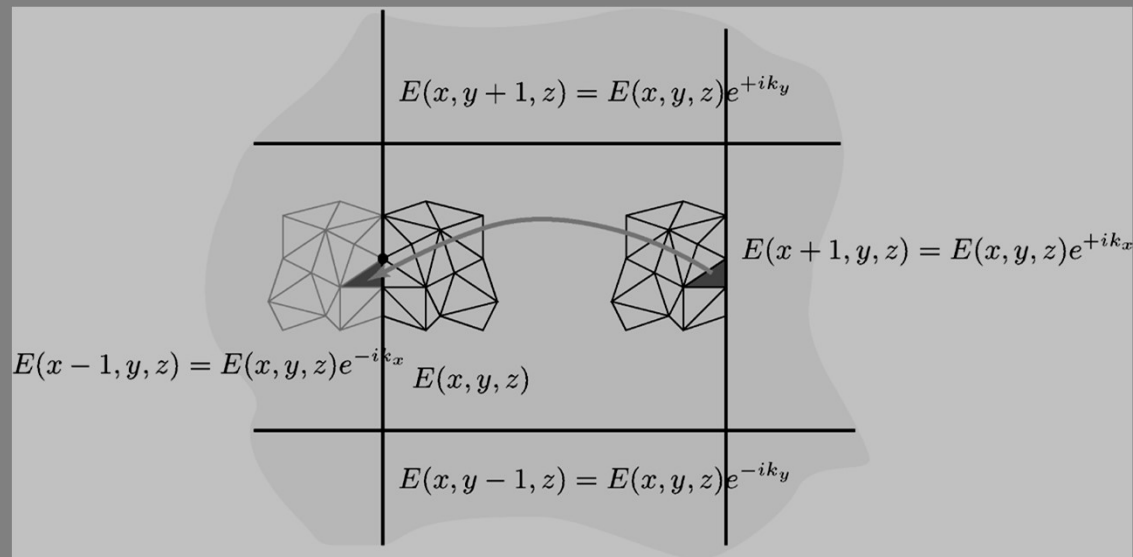
Bloch's theorem is sometimes stated in this alternative form:³ the eigenstates of H can be chosen so that associated with each ψ is a wave vector $\mathbf{k}$ such that

$$\psi(\mathbf{r} + \mathbf{R}) = e^{i\mathbf{k} \cdot \mathbf{R}} \psi(\mathbf{r})$$

'When I started to think about it, I felt that the main problem was to explain how the electrons could sneak by all the ions in a metal....

By straight Fourier analysis I found to my delight that the wave differed from the plane wave of free electrons only by a periodic modulation'

F. BLOCH



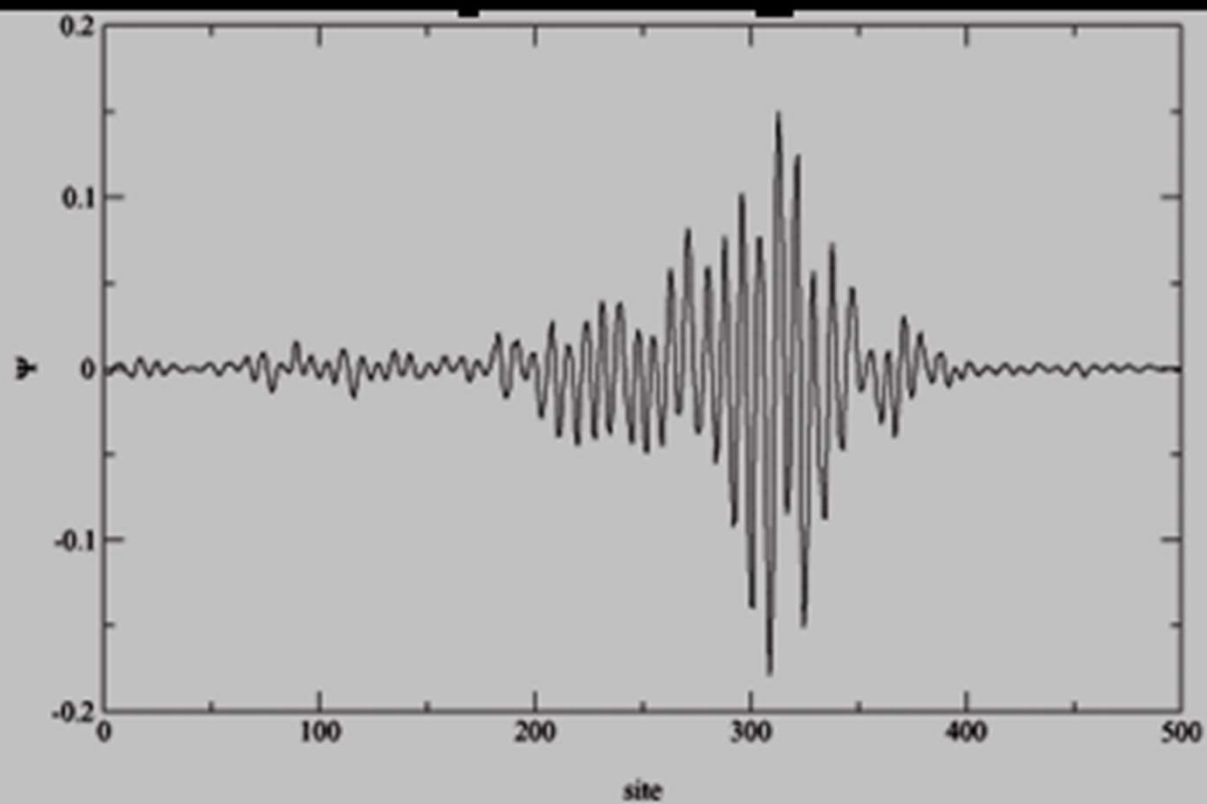


FIG. 4. Wave function of a delocalized state with energy close to 10.5 eV.

Proof of Bloch's Theorem

$$T_{\mathbf{R}}\psi(\mathbf{r}) = \psi(\mathbf{r} + \mathbf{R})$$

$$T_{\mathbf{R}}H(\mathbf{r})\psi(\mathbf{r}) = H(\mathbf{r}+\mathbf{R})\psi(\mathbf{r}+\mathbf{R}) = H(\mathbf{r})\psi(\mathbf{r}+\mathbf{R}) = H(\mathbf{r})T_{\mathbf{R}}\psi(\mathbf{r})$$

Translation operator commutes with Hamiltonian.

$$[H, T_{\mathbf{R}}] = 0$$

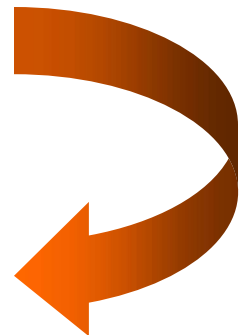
so they share the same eigenstates

$$H\psi(\mathbf{r}) = E\psi(\mathbf{r}) \qquad T_{\mathbf{R}}\psi(\mathbf{r}) = c(\mathbf{R})\psi(\mathbf{r})$$

$$T_{\mathbf{R}}T_{\mathbf{R}'}\psi(\mathbf{r}) = c(\mathbf{R})T_{\mathbf{R}'}\psi(\mathbf{r}) = c(\mathbf{R})c(\mathbf{R}')\psi(\mathbf{r})$$

$$T_{\mathbf{R}}T_{\mathbf{R}'}\psi(\mathbf{r}) = T_{\mathbf{R}+\mathbf{R}'}\psi(\mathbf{r}) = c(\mathbf{R} + \mathbf{R}')\psi(\mathbf{r})$$

$$c(\mathbf{R} + \mathbf{R}') = c(\mathbf{R})c(\mathbf{R}')$$



We can always write: $c(\mathbf{a}_i) = e^{2\pi i \mathbf{x}_i}$

Now let $\mathbf{a}_i$ be three primitive vectors for the Bravais lattice.

$$\mathbf{R} = n_1 \mathbf{a}_1 + n_2 \mathbf{a}_2 + n_3 \mathbf{a}_3$$

We had proved that:

$$c(\mathbf{R} + \mathbf{R}') = c(\mathbf{R})c(\mathbf{R}')$$

Thus: $C(\mathbf{R}) = C(n_1 \mathbf{a}_1 + n_2 \mathbf{a}_2 + n_3 \mathbf{a}_3) = C(n_1 \mathbf{a}_1)C(n_2 \mathbf{a}_2)C(n_3 \mathbf{a}_3)$

$$C(n_i \mathbf{a}_i) = C(\mathbf{a}_i)^{n_i}$$

Why ?

$$C(n_i \mathbf{a}_i) = C(\underbrace{\mathbf{a}_i + \mathbf{a}_i + \dots + \mathbf{a}_i}_{n_i}) = \underbrace{C(\mathbf{a}_i)C(\mathbf{a}_i)\dots C(\mathbf{a}_i)}_{n_i} = C(\mathbf{a}_i)^{n_i}$$



$$c(\mathbf{R}) = c(\mathbf{a}_1)^{n_1} c(\mathbf{a}_2)^{n_2} c(\mathbf{a}_3)^{n_3}$$

$$C(\mathbf{R}) = \exp\{2\pi i x_1\}^{n_1} \exp\{2\pi i x_2\}^{n_2} \exp\{2\pi i x_3\}^{n_3}$$

$$C(\mathbf{R}) = \exp\{2\pi i x_1\}^{n_1} \exp\{2\pi i x_2\}^{n_2} \exp\{2\pi i x_3\}^{n_3}$$

$$C(\mathbf{R}) = \exp\{2\pi i n_1 x_1\} \exp\{2\pi i n_2 x_2\} \exp\{2\pi i n_3 x_3\}$$

$$C(\mathbf{R}) = \exp\{2\pi i (n_1 x_1 + n_2 x_2 + n_3 x_3)\}$$

$$c(\mathbf{R}) = e^{i\mathbf{k} \cdot \mathbf{R}}$$

$$\mathbf{K} \cdot \mathbf{R} = 2\pi i (n_1 x_1 + n_2 x_2 + n_3 x_3)$$

Why ?

$$\mathbf{k} = x_1 \mathbf{b}_1 + x_2 \mathbf{b}_2 + x_3 \mathbf{b}_3 \quad \mathbf{R} = n_1 \mathbf{a}_1 + n_2 \mathbf{a}_2 + n_3 \mathbf{a}_3$$

$$\mathbf{b}_i \cdot \mathbf{a}_j = 2\pi \delta_{ij}$$

$$T_{\mathbf{R}} \psi = \psi(\mathbf{r} + \mathbf{R}) = c(\mathbf{R}) \psi(\mathbf{r}) = e^{i\mathbf{k} \cdot \mathbf{R}} \psi(\mathbf{r})$$



Plane Wave

Bloch's Wave

