

BAND STRUCTURE CALCULATION IN TIGHT BINDING MODEL

$$\psi_k = \frac{1}{\sqrt{N}} \sum_j e^{ik \cdot r_j} \phi(\mathbf{r} - \mathbf{r}_j)$$

Crystal electron wave function Atomic electron wave function

To calculate $E(k)$ we do as follows:

$$E(k) = \langle \psi_k | H | \psi_k \rangle = \int \psi_k^* H \psi_k dv$$

$$E(k) = \frac{1}{N} \sum_j \sum_m e^{ik \cdot (r_j - r_m)} \langle \phi(\mathbf{r} - \mathbf{r}_m) | H | \phi(\mathbf{r} - \mathbf{r}_j) \rangle$$

Just for simplicity we define that: $\{\phi_j = \phi(\mathbf{r} - \mathbf{r}_j), \phi_m = \phi(\mathbf{r} - \mathbf{r}_m)\}$

$$E(k) = \frac{1}{N} \sum_j \sum_m e^{ik \cdot (r_j - r_m)} \langle \phi_m | H | \phi_j \rangle$$

We calculate the above summation for an special atom located at $r_j = 0$, and then multiply it by N

$$r_j = 0, \quad \rho_m = r_m - r_j = r_m$$

$$E(k) = N \times \frac{1}{N} \sum_m e^{-ik \cdot \rho_m} \int \phi^*(\mathbf{r} - \rho_m) H \phi(\mathbf{r}) dv$$

Summation is over all the atoms with respect to the atom located at the origin, i.e. $r_j = 0$

Okay, now if we consider only nearest neighbor atoms and neglect the overlap with the other atomic orbitals:

For the main atom at the origin $\rho_m = 0$ $\int \phi^*(\mathbf{r}) H \phi(\mathbf{r}) dv = -\alpha$

$\int \phi^*(\mathbf{r} - \rho) H \phi(\mathbf{r}) dv = -\gamma$ For the nearest neighbors $\rho_m = \rho$

$$E(k) = \sum_m e^{-ik \cdot \rho_m} \int \phi^*(\mathbf{r} - \rho_m) H \phi(\mathbf{r}) dv$$

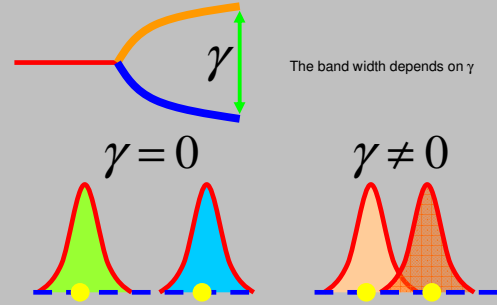
$$E(k) = -\alpha - \gamma \sum_m e^{-ik \cdot \rho_m}$$

This term originates from the interaction of the main atom with its nearest neighbors

For the isolated atom, i.e. localized electrons

$$\int \phi^*(\mathbf{r} - \rho) H \phi(\mathbf{r}) dv = -\gamma$$

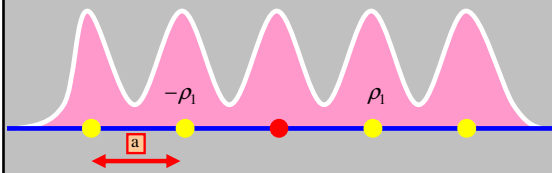
γ is called the overlap integral. γ is a parameter that causes to be changed the energy levels to energy bands.



$$E(k) = -\alpha - \gamma \sum_m e^{-ik \cdot \rho_m}$$

Over nearest neighbors

For a 1D lattice:



$$E(k) = -\alpha - \gamma \sum_{m=-1}^1 e^{-ik \cdot \rho_m} = -\alpha - \gamma (e^{-ik \cdot \rho_{-1}} + e^{-ik \cdot \rho_1})$$

$$= -\alpha - \gamma (e^{ika} + e^{-ika}) = -\alpha - \gamma (2 \cos(ka))$$

$$E(k) = -\alpha - \gamma (2 \cos(ka))$$

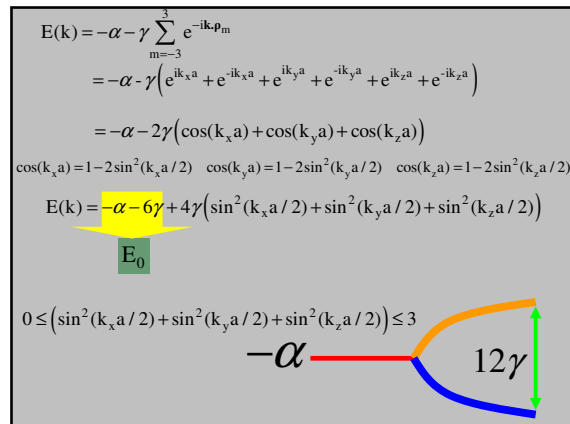
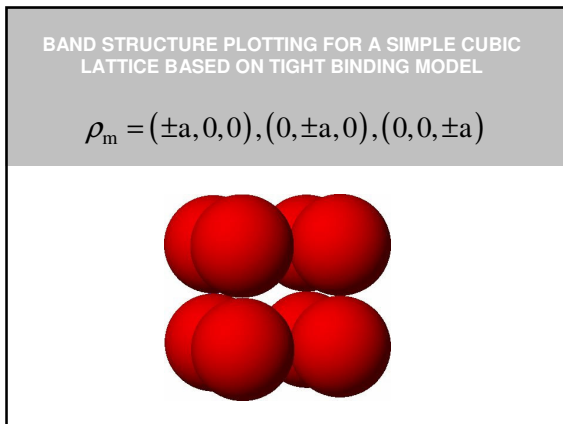
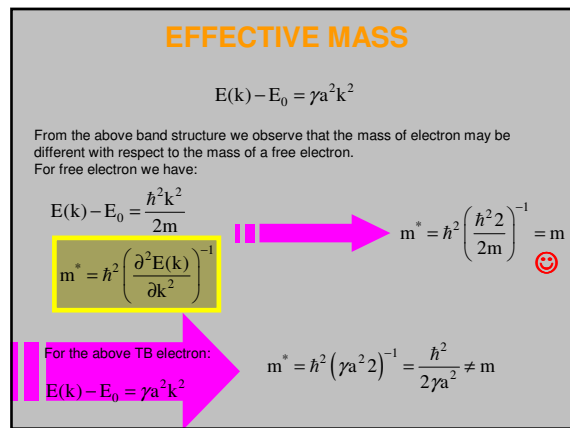
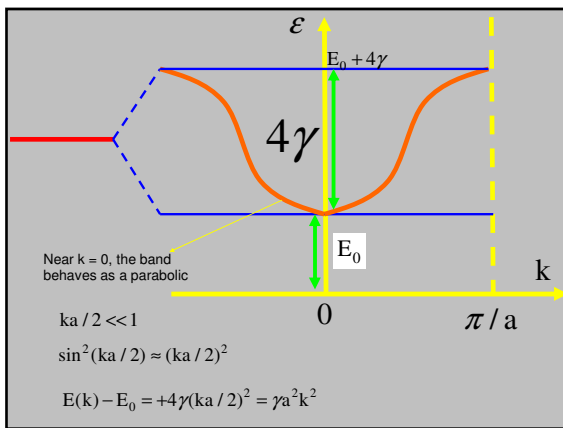
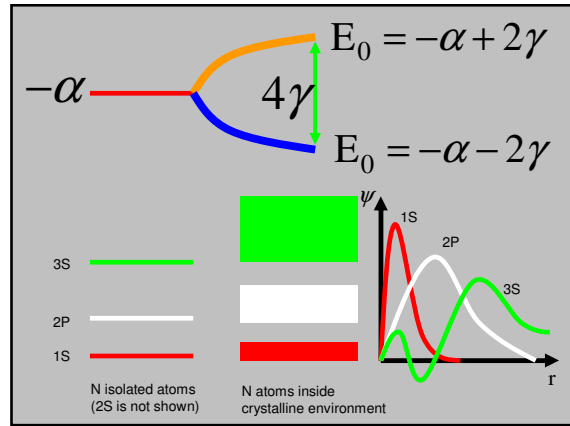
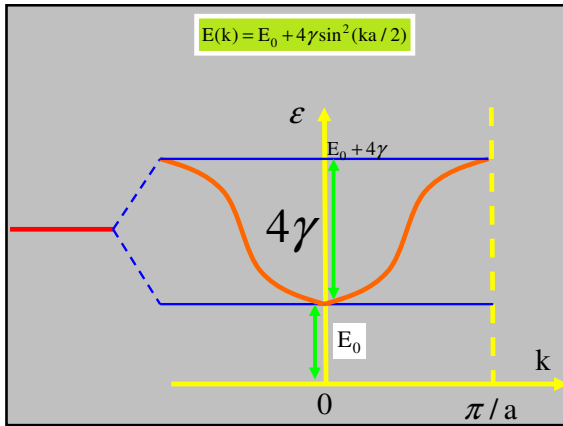
$$\cos(ka) = 1 - 2 \sin^2(ka/2)$$

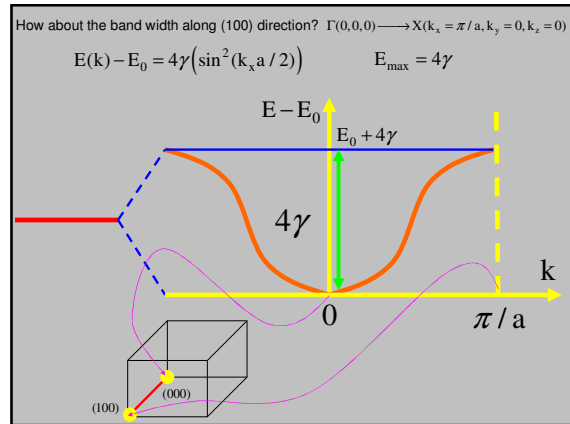
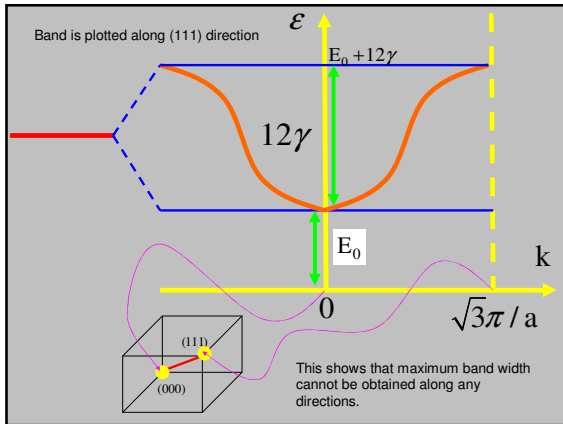
$$E(k) = -\alpha - \gamma (2 - 4 \sin^2(ka/2))$$

$$E(k) = -\alpha - 2\gamma + 4\gamma \sin^2(ka/2)$$

E_0

$$E(k) = E_0 + 4\gamma \sin^2(ka/2)$$





Near the Γ point the band structure is spherical.

$(\sin(k_x a/2) \approx k_x a/2, \sin(k_y a/2) \approx k_y a/2, \sin(k_z a/2) \approx k_z a/2)$

$E(k) = E_0 + 4\gamma(\sin^2(k_x a/2) + \sin^2(k_y a/2) + \sin^2(k_z a/2))$

$E(k) - E_0 \approx 4\gamma((k_x a/2)^2 + (k_y a/2)^2 + (k_z a/2)^2)$

$\approx \gamma a^2(k_x^2 + k_y^2 + k_z^2) = \gamma a^2 k^2$

To find the effective mass: $m^* = \hbar^2(\gamma a^2)^{-1} = \frac{\hbar^2}{2\gamma a^2} \neq m$

The band is deformed near the edge of 1BZ so that it be perpendicular to the edge:

$E(k) = E_0 + 4\gamma(\sin^2(k_x a/2) + \sin^2(k_y a/2) + \sin^2(k_z a/2))$

$\frac{\partial E(k)}{\partial k} \bigg|_{k_x = \frac{\pi}{a}} = 4\gamma(2 \times a/2 \cos(k_x a/2) \sin(k_x a/2)) \bigg|_{k_x = \frac{\pi}{a}} = 0$

Is this the case for free electron?

ENERGY CONTOUR IN K-SPACE

In order to plot energy contour, we would first plot the Fermi surface, and then find intersections of an specific plane with the Fermi surface:

$E(k) = -\alpha - 6\gamma + 4\gamma(\sin^2(k_x a/2) + \sin^2(k_y a/2) + \sin^2(k_z a/2))$

$= \text{const} \tan t$

$\sin^2(k_x a/2) + \sin^2(k_y a/2) + \sin^2(k_z a/2) = \text{const} \tan t$

PRACTICE

Plot the above Fermi surface in 3 wave vector space dimensional space.

The result of the last practice will be similar to the following Fermi surface:

Near the Γ -point the shape of the Fermi surface is spherical. But far from it the deviation is observed.

A typical energy contour can be as follows:

k_z

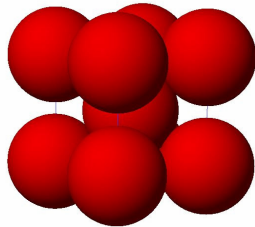
k_y

k_x

Γ -point is here

BAND STRUCTURE PLOTTING FOR A BODY CENTERED CUBIC LATTICE BASED ON TIGHT BINDING MODEL

$$\rho_m = (\pm a/2, \pm a/2, \pm a/2)$$

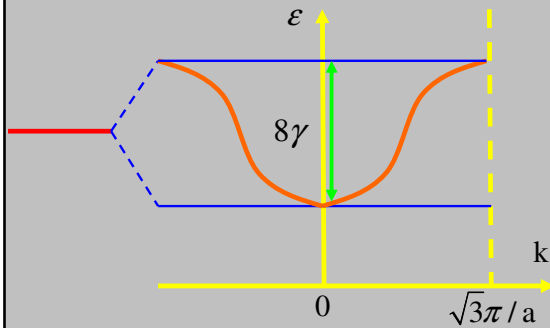


$$E(k) = -\alpha - \gamma \sum_{m=1}^3 e^{-ik \cdot \rho_m} \\ = -\alpha - \gamma \left\{ e^{-i(k_x a/2 + k_y a/2 + k_z a/2)} + e^{-i(k_x a/2 + k_y a/2 - k_z a/2)} + \right. \\ \left. e^{-i(k_x a/2 - k_y a/2 + k_z a/2)} + e^{-i(k_x a/2 - k_y a/2 - k_z a/2)} + \right. \\ \left. e^{-i(k_x a/2 + k_y a/2 - k_z a/2)} + e^{-i(k_x a/2 - k_y a/2 + k_z a/2)} \right\}$$

$$E(k) = -\alpha - \gamma \left\{ 2 \cos(k_x a/2 + k_y a/2 + k_z a/2) + \right. \\ \left. 2 \cos(k_x a/2 - k_y a/2 + k_z a/2) + \right. \\ \left. 2 \cos(k_x a/2 + k_y a/2 - k_z a/2) \right\}$$

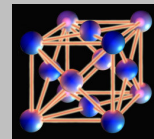
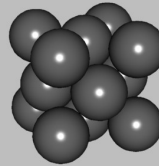
$$E(k) = -\alpha - 8\gamma \cos(k_x a/2) \cos(k_y a/2) \cos(k_z a/2)$$

Band is plotted along (111) direction



PRACTICE

Find the band structure of a fcc lattice based on the tight binding model:



$$E(k) = -\alpha - 4\gamma \left(\cos(k_x a/2) \cos(k_y a/2) + \cos(k_x a/2) \cos(k_z a/2) + \cos(k_y a/2) \cos(k_z a/2) \right)$$

PRACTICE

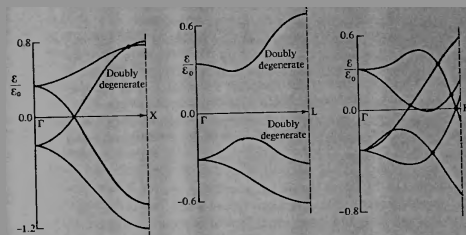


Figure 10.6

A tight-binding calculation of the 3d bands of nickel. (G. C. Fletcher, *Proc. Phys. Soc.*, A65, 192 (1952).) Energies are given in units of $\epsilon_0 = 1.349$ eV, so the bands are about 2.7 volts wide. The lines along which E is plotted are shown in Figure 10.5. Note the characteristic degeneracies along ΓX and ΓL , and the absence of degeneracy along ΓK . The great width of the bands indicates the inadequacy of so elementary a treatment.